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REACTIONS OF UO(+) WITH ATMOSPHERIC GASES.(U)

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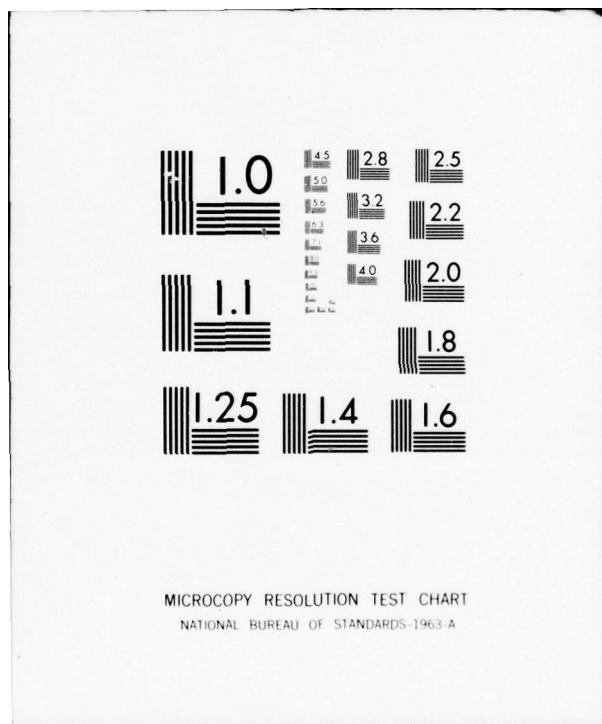
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REACTIONS OF  $UO_2^+$  WITH ATMOSPHERIC GASES

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20. ABSTRACT (Continue on reverse side if necessary and identify by block number) The rate coefficient for the ion-molecule reaction of uranium monoxide ions with molecular oxygen at thermal energies has been measured using a "magnetic bottle" apparatus. The UO <sub>2</sub> <sup>+</sup> ions were produced by the associative ionization reaction between U and N <sub>2</sub> O. The primary (UO <sub>2</sub> <sup>+</sup> ) and secondary (UO <sub>2</sub> <sup>+</sup> ) ions were selected by a quadrupole mass spectrometer and their variations vs. O <sub>2</sub> pressure were plotted on a X-Y recorder. The rate coefficient was determined to be 2.03 + 0.44 x 10 <sup>-9</sup> cm <sup>3</sup> /sec. <i>TEN TO THE MINUS 9TH POWER CUBIC CM</i>		

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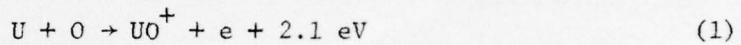
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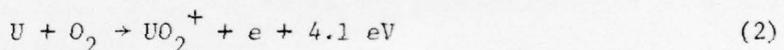
## REACTIONS OF $\text{UO}^+$ WITH ATMOSPHERIC GASES

### I. Introduction

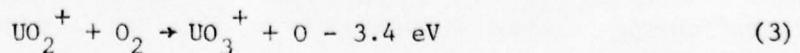
When uranium atoms are released in the upper atmosphere where atomic and molecular oxygen are present, the associative ionization reactions



and

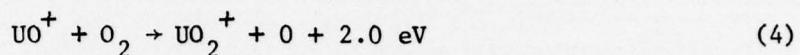


occur. The plasma produced by these reactions cannot be neutralized by the normal dissociative recombination reactions because of the endoergicity of the reactions inverse to Reactions (1) and (2), and hence a long-lived plasma which can cause interference to radar communications can be produced. Further, because of the exoergicity of Reactions (1) and (2), it is expected that both the  $\text{UO}^+$  and  $\text{UO}_2^+$  ions will be formed in excited states and subsequently radiate causing interference with optical and infra-red communication systems. There is firm experimental evidence from this laboratory that of the order of 1% of the  $\text{UO}_2^+$  is formed with internal energies higher than 3.4 eV from the observation that the reaction



proceeds when the  $\text{UO}_2^+$  is formed by Reaction (2).

It is of interest to determine the rate at which the reaction



proceeds. Where there are both O and O<sub>2</sub> present in the atmosphere this reaction can reduce the UO<sup>+</sup> spectrum with time and enhance the UO<sub>2</sub><sup>+</sup> spectrum. The exoergicity of Reaction (4) may also cause an alteration of the UO<sub>2</sub><sup>+</sup> spectrum, through population of excited states in addition to those populated by Reaction (2).

This report describes the experimental measurement of the rate coefficient for Reaction (4) at thermal energies.

### II. Experimental Approach

The approach employs a "magnetic bottle" apparatus<sup>1</sup> which is shown in Fig. 1. Within the vacuum system are coils which produce a magnetic field bounded by two magnetic mirrors. A beam of uranium atoms crosses the magnetic field at the center of the bottle and ions are formed along the atomic beam's length by associative ionization reactions with gases admitted at very low pressure into the vacuum chamber. Those ions that are formed within the magnetic bottle are trapped in the bottle, moving back and forth between the mirrors on helical trajectories along the magnetic field lines. Eventually elastic collisions of the ions with the added gas change the pitch angle of the ions to a sufficiently low value that the ions can escape through the mirrors bounding the bottle. A quadrupole mass filter is placed outside one of the mirrors and detects the escaping ions.

The magnetic bottle used was produced by a set of ten water-cooled coils, each being 5.1 cm i.d., 15.2 cm o.d., and 2 cm thick. A gold-plated

brass cylinder of 5 cm diameter was inserted in the column of the coils and maintained at -3V. The bottle provided a field strength of 1500 G at the center and 2500 G at the mirrors.

The uranium atom beam was produced in a source consisting of a tungsten tube of 5 cm length and 6 mm diameter, made by rolling a 5.7 cm x 10.2 cm tungsten foil 0.025 mm thick. A hole of 1.6 mm diameter was placed in the side wall of the tube at the center along its length although we have found in previous experiments that the hole is actually unnecessary due to the ability of U atoms to diffuse through the hot tungsten foils. The hole was made for the convenience of recharging U supply into the W tube. Two pieces of 1.6-mm diameter U rod, each 5 mm long, were confined in the central portion of the tube by two waddings of Ta fine wires. The tube was fitted into molybdenum caps at each end which in turn were bolted onto a water-cooled copper jacket which both carried the electrical current and dissipated furnace heat. A current of about 200 A at 2.5 V heated the central portion of the tube to about 2100°K.

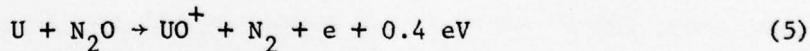
The U atom beam was formed by passing through two collimators attached to the copper cooling jacket before entering the reaction region through a 4 mm x 8 mm entrance aperture in the side of the gold-plated brass cylinder. Proper biasings (the furnace tube at +6V, the first collimator at the ground potential, and the second at +70V) prevented both ions and electrons produced in the U atom source from reaching the reaction region.

A pair of parallel plates situated on the other side of the magnetic coil column and across from the U atom source was used to monitor  $\text{UO}_2^+$  formed by Reaction (2) when the U beam intersected a known amount of  $\text{O}_2$  admitted into the vacuum. The  $\text{UO}_2^+$  ion current measured by the condenser method at a fixed

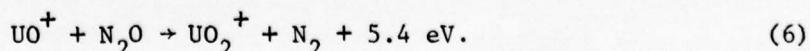
$O_2$  pressure indicates the U atom beam current and is important in the measurements because the U atom beam current tends to gradually decrease with time.

The primary and secondary ions were extracted through a 9.9-mm exit aperture at one end of the magnetic bottle, some 16 cm from the reaction region, but an extractor normally biased at -2V. The ions were then m/e selected by a quadrupole mass spectrometer and detected by an EMI 21-state venetian blind electron multiplier. The ion current output from the multiplier was measured by an Extranuclear electrometer and the output of the latter was fed into the y-axis of a X-Y recorder. The pressures of the reactant gases were monitored on the x-axis of the X-Y recorder. The pressure readings of  $N_2O$  and  $O_2$  added to the vacuum were furnished by a Hughes type 6578 ionization gauge which was calibrated for  $N_2O$  and  $O_2$  against a Consolidated GM-110 McLeod gauge.

In the present experiments we added two gases,  $N_2O$  and  $O_2$ , in the vacuum. The primary ions were  $UO^+$  and  $UO_2^+$  formed by the reaction



and Reaction (2). The respective rate coefficients,  $k_5$  and  $k_2$ , are known<sup>2,3</sup> to be  $1.51 \times 10^{-13} \text{ cm}^3/\text{sec}$  and  $1.04 \times 10^{-12} \text{ cm}^3/\text{sec}$ .  $UO_2^+$  can be produced additionally by the reaction



The rate coefficient,  $k_6$ , has also been determined to be  $2.73 \times 10^{-10} \text{ cm}^3/\text{sec}$  by the authors in a previous magnetic bottle experiment. If the rate coeffi-

cient for the reaction of interest, Reaction (4), is designated as  $k_4$ , we can write down the rate of change of the concentrations  $n_1$  (for  $\text{UO}_2^+$ ) and  $n_2$  (for  $\text{UO}_2^+$ ) as

$$\frac{dn_1}{dt} = k_5 n_u n_a - (k_6 n_a + k_4 n_b) n_1 - L n_1 \quad (7)$$

and

$$\frac{dn_2}{dt} = k_2 n_u n_b + (k_6 n_a + k_4 n_b) n_1 - L n_2 \quad (8)$$

where  $n_u$  is the concentration in the U atom beam, and  $n_a$  and  $n_b$  are the number densities of  $\text{N}_2\text{O}$  and  $\text{O}_2$ , respectively, admitted into the vacuum.  $L$  is the loss rate due to escape of the ions from the mirrors, which is proportional to  $n_a + n_b$ .

In the steady state, Equations (7) and (8) solve to give the ratio of  $\text{UO}_2^+$  and  $\text{UO}_2^+$  as

$$\frac{n_2}{n_1} = \frac{1}{L} [k_6 n_a + (k_4 + \frac{k_2 k_6}{k_5} + \frac{k_2}{k_5 n_a} \cdot L) n_b + \frac{k_2 k_4}{k_5 n_a} n_b^2]. \quad (9)$$

Since  $k_6$  is known and  $n_a$  is measured in the experiment, the loss coefficient,  $L_o$ , is determined for  $n_b = 0$  from the ratio of the secondary to primary ions when only the  $\text{N}_2\text{O}$  is present. The loss coefficient with added  $\text{O}_2$  is then given by  $L = L_o(1 + (n_b/n_a))$  and this can be used to obtain  $k_4$  from the linear portion of a curve of  $n_2/n_1$ , for fixed  $n_a$ , as a function of  $n_b$ , i.e., the slope as the added oxygen gas pressure approaches zero.

Because of cumulative uncertainties in the linear coefficient expression, however, it is more satisfactory to use the coefficient of the quadratic term in Equation (9) to obtain  $k_4$ .

### III. Method

In the present experiment, the  $\text{UO}^+$  ions that were produced by Reaction (5) at a fixed  $\text{N}_2\text{O}$  pressure reacted with  $\text{O}_2$ , which was added to the vacuum chamber in addition to the fixed  $\text{N}_2\text{O}$  pressure, forming  $\text{UO}_2^+$  ions that were also produced simultaneously by Reactions (2) and (6). Figure 2 shows the typical curves of  $\text{UO}_2^+$  and  $\text{UO}^+$  ion currents vs.  $\text{O}_2$  pressure obtained on a X-Y recorder for a fixed  $\text{N}_2\text{O}$  pressure. The pressure of  $\text{N}_2\text{O}$  was typically  $3 \times 10^{-5}$  Torr while the pressure range of  $\text{O}_2$  was from 0 to  $1.5 \times 10^{-4}$  Torr. The measurements were repeated for a number of  $\text{N}_2\text{O}$  pressures.

Using data from these two curves, a plot of  $n_2/n_1$  (ratio of  $\text{UO}_2^+$  to  $\text{UO}^+$  currents) as a function of the number density of  $\text{O}_2$ ,  $n_b$ , was drawn for each fixed number density of  $\text{N}_2\text{O}$ ,  $n_b$ . The curve of  $n_2/n_1$  vs.  $n_b$  is a parabola, as expected by Equation (9). The linear term of Equation (9) can be obtained from the tangent to this curve at  $n_b = 0$ . Next, the square root of the difference for  $n_2/n_1$  between the parabola and the tangent was plotted as a function of  $n_b$ , forming a straight line. The slope,  $S$ , of this straight line was used to determine the rate coefficient of Reaction (4) using the relation

$$k_4 = \frac{k_5}{k_2} n_a L S^2. \quad (10)$$

#### IV. Results and Discussion

The rate coefficient for the ion-molecule reaction of  $\text{UO}^+ + \text{O}_2 \rightarrow \text{UO}_2^+ + \text{O}$  at thermal energies was found to be  $2.03 \pm 0.44 \times 10^{-9} \text{ cm}^3/\text{sec.}$

Under different experimental conditions, such as changes in the uranium atom source temperature (ranging from 2050 to 2200°K) and variations in the biasing of the reaction region (at -3V and +0.6 V), no obvious difference was observed for the rate coefficient within the experimental errors.

At thermal energies, it appears that the ion-molecule reaction of  $\text{UO}^+ + \text{O}_2 \rightarrow \text{UO}_2^+ + \text{O}$  is almost one order of magnitude faster than that of  $\text{UO}^+ + \text{N}_2\text{O} \rightarrow \text{UO}_2^+ + \text{N}_2$ , the rate coefficient of the latter being  $2.73 \times 10^{-10} \text{ cm}^3/\text{sec.}$

Since the loss coefficient, L, is related to the number density of oxygen,  $n_b$ , by  $L_o(1 + \frac{n_b}{n_a})$ , Equation (10) indicates that the rate coefficient,  $k_4$ , for  $\text{UO}^+ + \text{O}_2 \rightarrow \text{UO}_2^+ + \text{O}$  is not independent of  $n_b$ . Hence the rate coefficient of  $2.03 \times 10^{-9} \text{ cm}^3/\text{sec}$  may be considered as the lower limit of the reaction investigated.

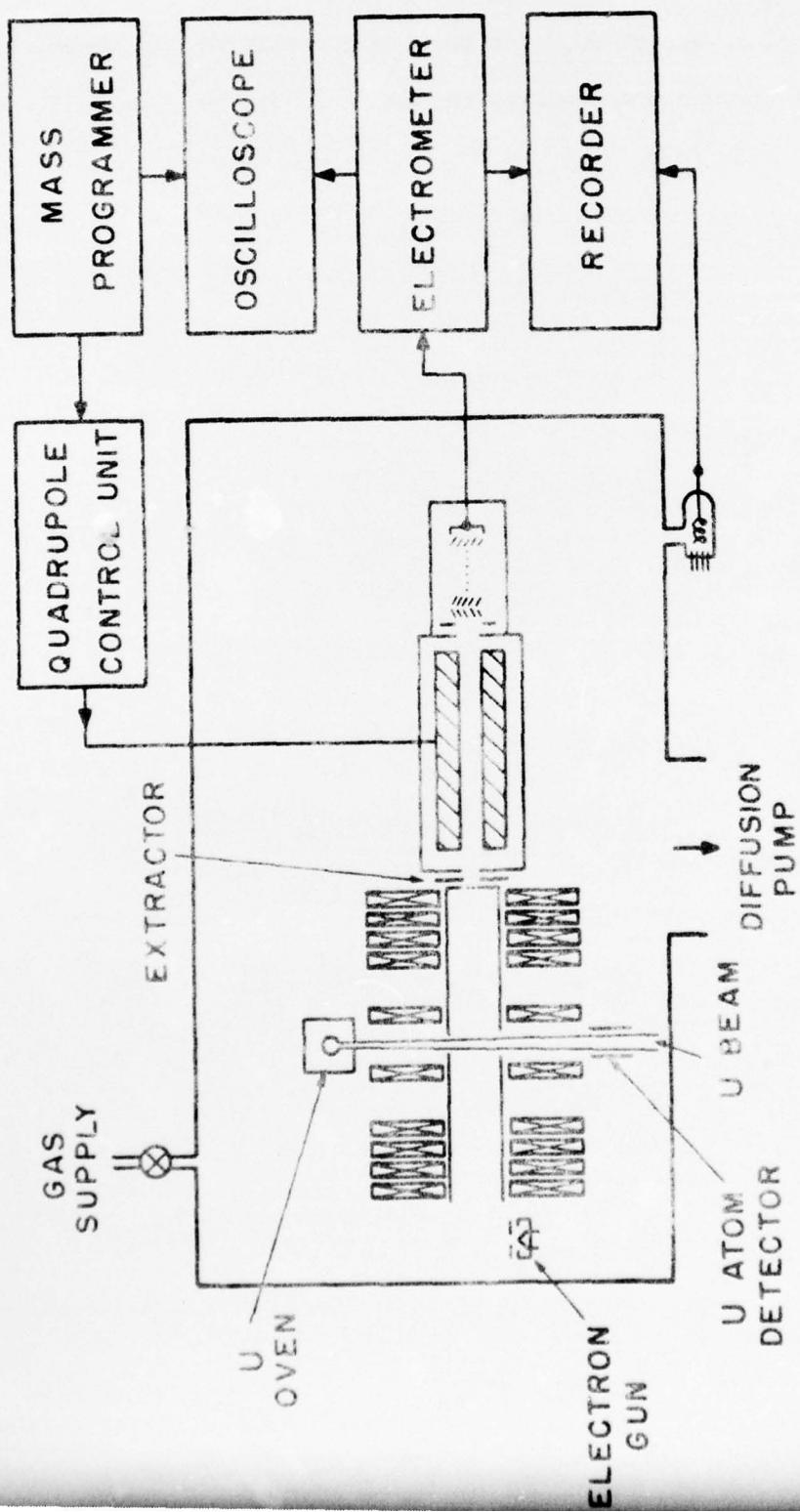
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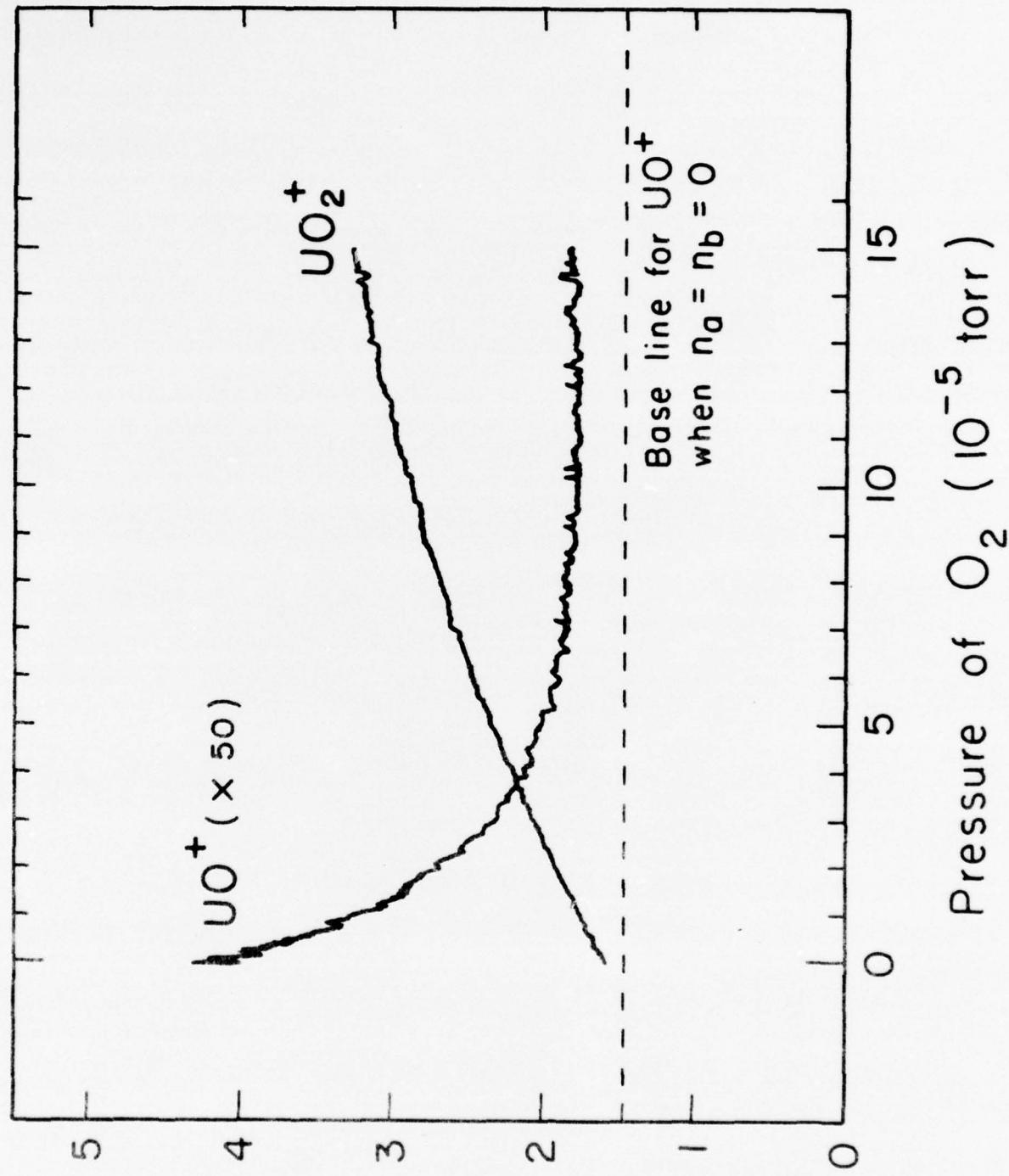
Figure Captions

Figure 1. "Magnetic Bottle" apparatus.

Figure 2. Typical curves of  $UO_2^+$  and  $UO^+$  ion currents vs.  $O_2$  pressure  
for a fixed number density of  $N_2O$  at  $1.07 \times 10^{12} \text{ cm}^{-3}$ .



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